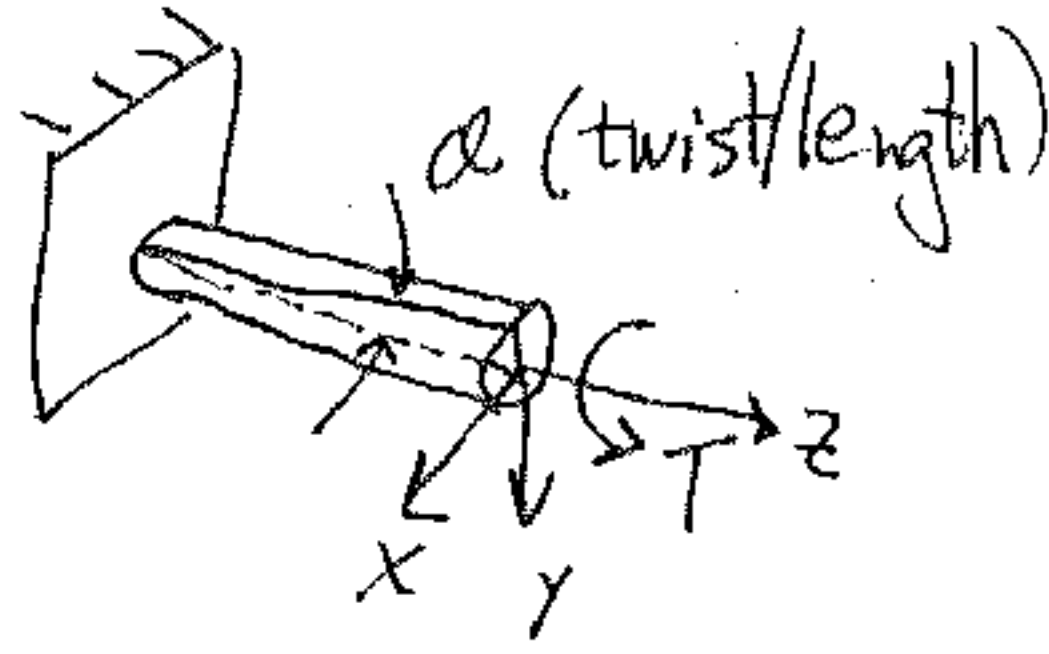


Torsions St. Venant's semi-inverse method

circular bar



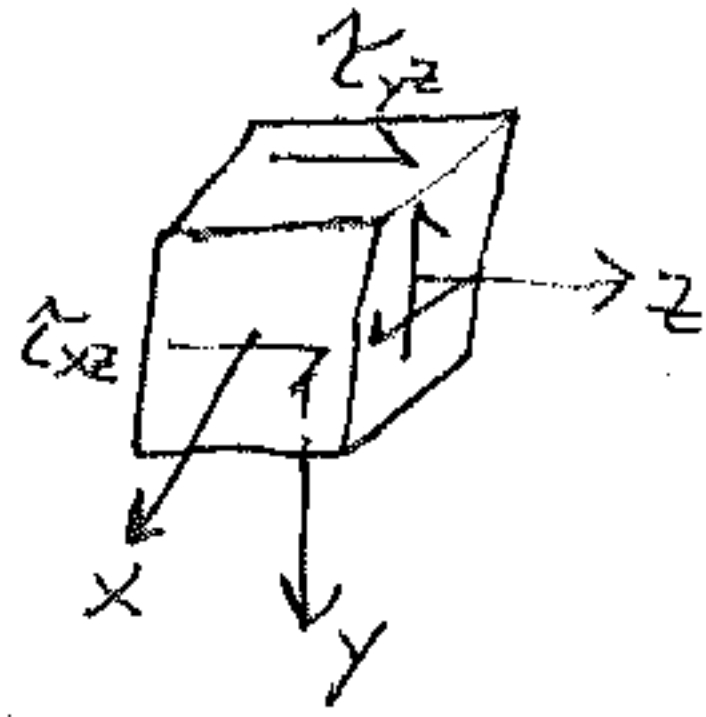
torsion rigidity  $C = \frac{T}{\alpha}$

assume: plane sections remain plane

$$u = u_x = -y\alpha z$$

$$v = u_y = \alpha z x$$

$$w = 0 \text{ (plane sections remain plane)}$$



$$\epsilon_{ij} = \frac{1}{2}(u_{i,j} + u_{j,i})$$

$$\epsilon_{xx} = \epsilon_{yy} = \epsilon_{zz} = 0$$

$$\begin{aligned} \epsilon_{xz} &= \frac{1}{2}(u_{x,z} + u_{z,x}) \\ &= \frac{1}{2}(-\alpha) \end{aligned}$$

$$\epsilon_{yz} = \frac{1}{2}(u_{y,z} + u_{z,y}) = \frac{1}{2}\alpha x$$

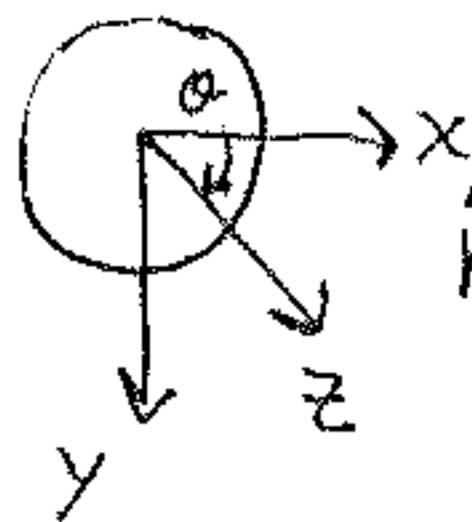
$$\tau = G\gamma$$

$$\tau_{xz} = -G\alpha y$$

$$\tau_{yz} = G\alpha x$$

on surface away from ends

$$\underline{t} = 0 = t_x = t_y = t_z$$



$$\hat{n} = n_x \hat{i} + n_y \hat{j} = \frac{x}{r} \hat{i} + \frac{y}{r} \hat{j}$$

can use  $\cos\alpha, \sin\alpha$  also

$$t_j = \sigma_{ji} n_j$$

$$\frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial x^2} = GQ \left( \frac{\partial^2 \phi}{\partial x \partial y} - 1 \right) - GQ \left( \frac{\partial^2 \phi}{\partial y \partial x} + 1 \right)$$

$$= \underbrace{-2GQ}_F$$

Boundary conditions

on the surface,  $z=0$

$$\tau_{xz} = \tau_{yz} = 0$$

$$\tau_{xz} n_x + \tau_{yz} n_y = 0$$

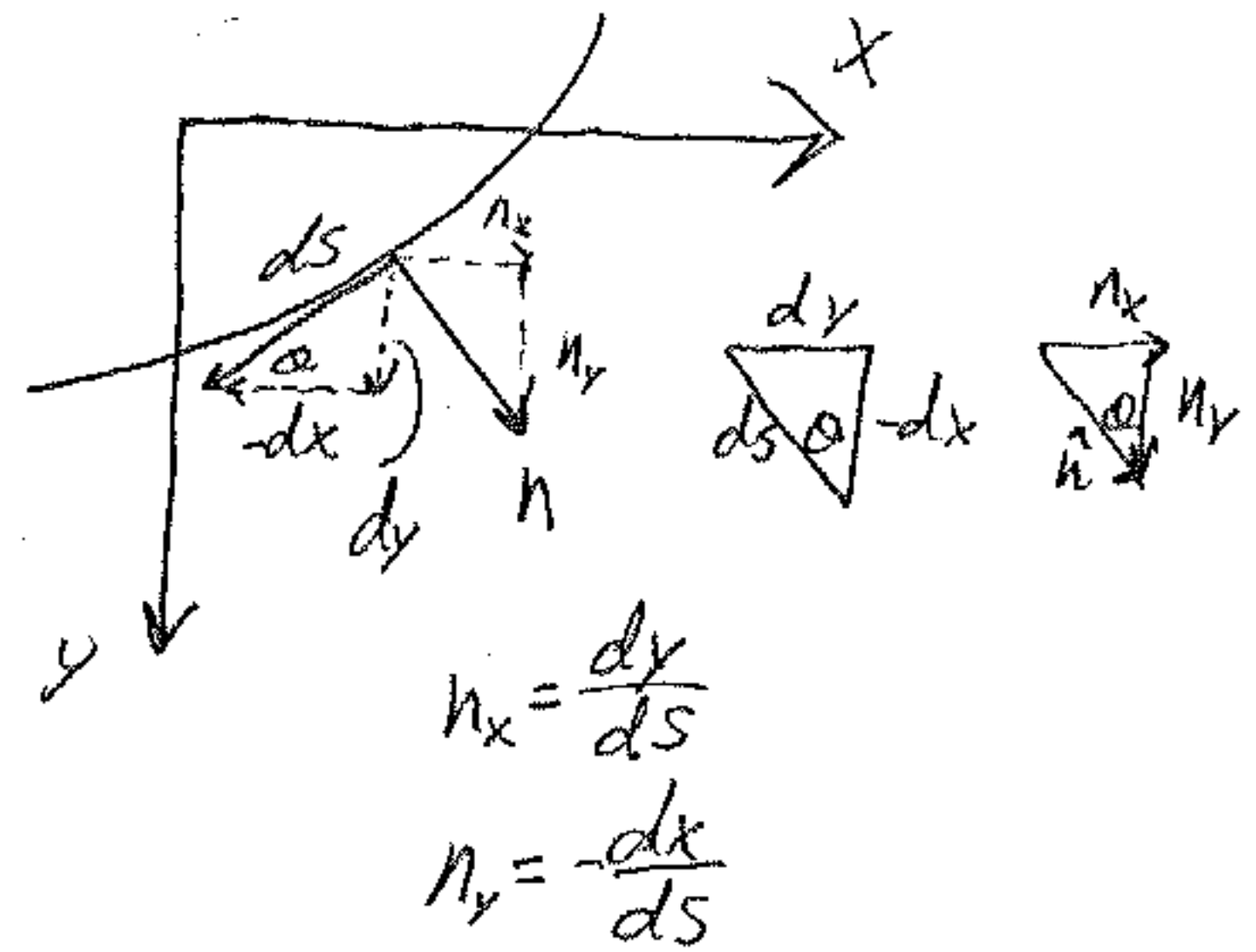
$$\frac{\partial \phi}{\partial y} n_x - \frac{\partial \phi}{\partial x} n_y = 0$$

$$\frac{\partial \phi}{\partial y} \frac{\partial y}{\partial s} + \frac{\partial \phi}{\partial x} \frac{\partial x}{\partial s} = 0 = \frac{d\phi}{ds}$$

$\phi$  is constant on  $S$

$\phi = 0$  picked arbitrarily, see 7.6 section in book

$\tau_{xz} = \frac{\partial \phi}{\partial y} \therefore$  it doesn't matter if  $\phi = \text{const. or zero}$



BC. of bar on end

$$\hat{n} = \hat{i} = \hat{k}$$

$$n_z = 1$$

$$t_x = \tau_{xz}$$

$$t_y = \tau_{yz}$$

$$\Sigma F_x = 0 = \int_A \tau_{xz} dA = \int_A \frac{\partial \phi}{\partial y} dx dy = \int \phi \Big|_{y_1}^{y_2} dx$$

$$\Sigma M = M_t, \quad \Sigma F_y = \int_A -\frac{\partial \phi}{\partial x} dx dy = -\int \phi \Big|_{x_1}^{x_2} dy = 0$$

$$M_t = \int_A (t_y x - t_x y) dx dy$$

$$= \int_A -\frac{\partial \phi}{\partial x} x dx dy + \int_A -\frac{\partial \phi}{\partial y} y dx dy$$

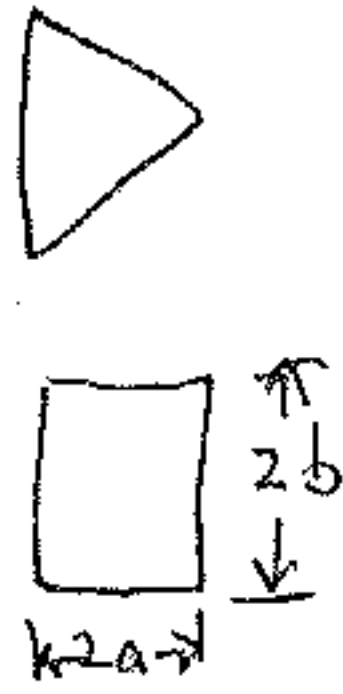
integrate by parts with  $\phi = 0$  on boundary

giving  $M_t = 2 \int_A \phi dx dy$

53

Elliptical cross-sections fairly easy

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$



these are harder

try  $\phi = m \left( \frac{x^2}{a^2} + \frac{y^2}{b^2} - 1 \right) = 0$

recall  $\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = -2GQ = F$

$$F = m \left( \frac{2}{a^2} + \frac{2}{b^2} \right)$$

$$m = \frac{F}{2} \left( \frac{a^2 b^2}{a^2 + b^2} \right)$$

$$\phi = \frac{F}{2} \frac{a^2 b^2}{a^2 + b^2} \left( \frac{x^2}{a^2} + \frac{y^2}{b^2} - 1 \right) = 0$$

now use,

$$M_t = 2 \int_A \phi \, dx \, dy$$

$$= \frac{a^2 b^2 F}{a^2 + b^2} \left[ \underbrace{\int_A \frac{x^2}{a^2} \, dx \, dy}_{\frac{1}{a^2} I_y} + \underbrace{\int_A \frac{y^2}{b^2} \, dx \, dy}_{\frac{1}{b^2} I_x} - \underbrace{\int_A dx \, dy}_A \right]$$

$$\frac{\pi}{4} \frac{a^3 b}{a^2} + \frac{\pi}{4} \frac{b^3 a}{b^2} - \pi ab = -\frac{\pi}{2} ab$$

$$M_t = -\frac{\pi a^3 b^3 F}{2(a^2 + b^2)}$$

$$\phi = -\frac{M_t}{\pi ab} \left( \frac{x^2}{a^2} + \frac{y^2}{b^2} - 1 \right)$$

$$\tau_{xz} = \frac{-2 M_t y}{\pi a b^3}$$

$$\tau_{yz} = \frac{2 M_t x}{\pi a^3 b}$$

must work back using  $w = Q\psi(x, y)$

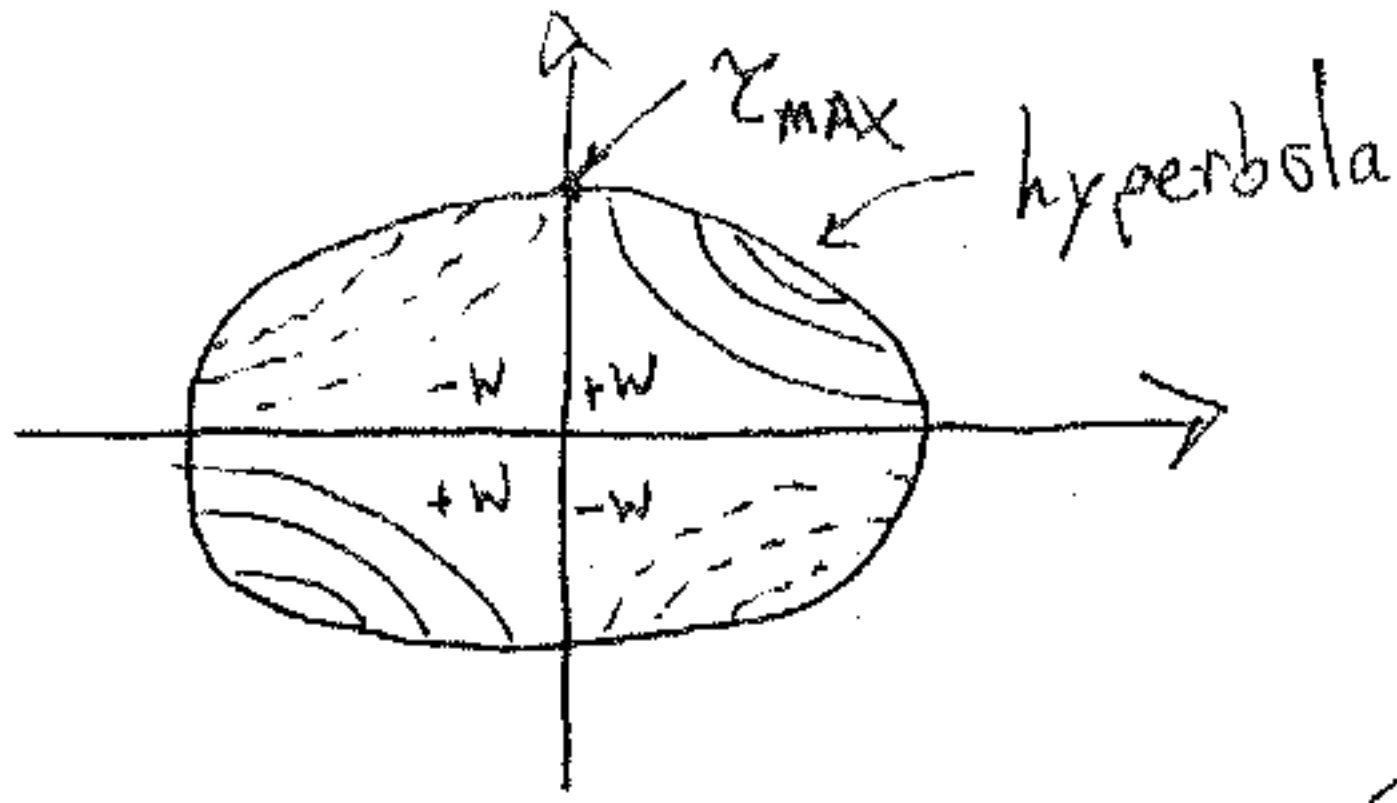
$$w = \frac{M_t (b^2 - a^2)}{\pi a^3 b^3 G} xy$$

$$Q = \frac{M_t (a^2 + b^2)}{\pi a^3 b^3 G}$$

Back substitution using equilibrium and shear stress

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = -2GQ = F$$

50



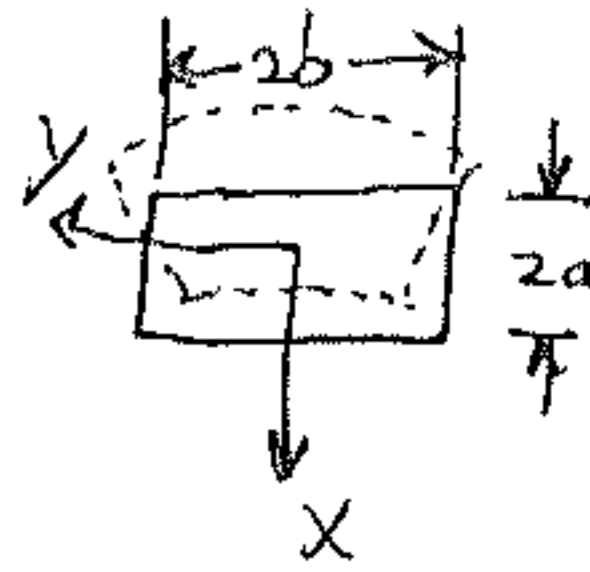
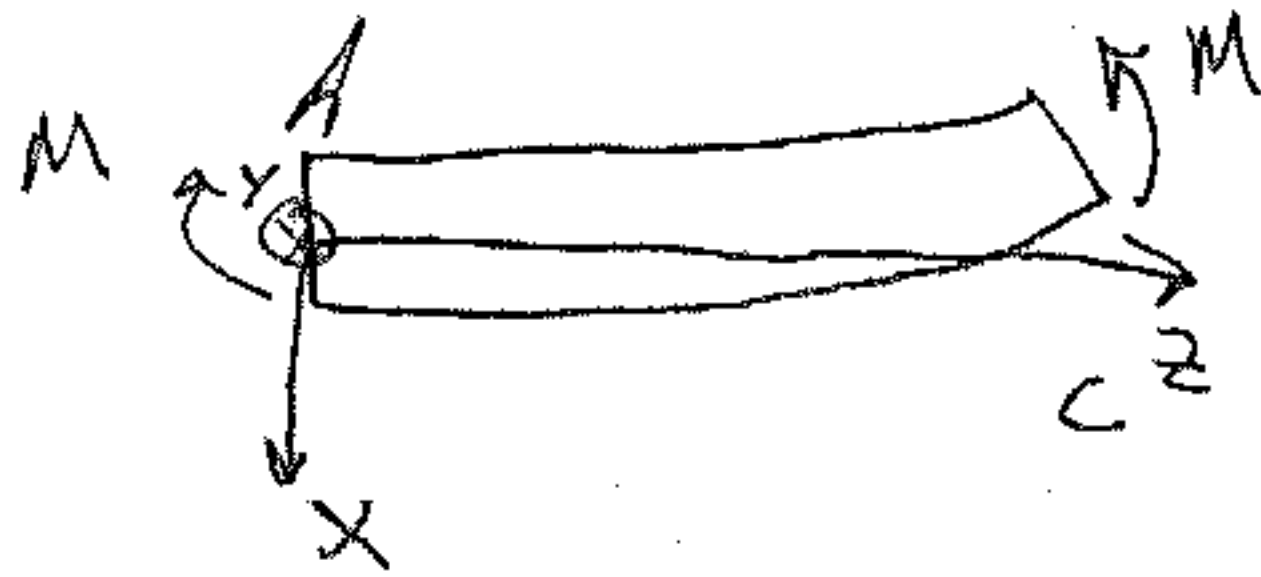
Torsional rigidity

$$C = \frac{I}{\theta} = \frac{G}{4\pi^2} \frac{A^4}{I_p}$$

A-area =  $\pi ab$   
 $I_p = I_x + I_y$

$$\begin{aligned} \tau_{max} &= \text{on the boundary at } y=b \text{ (minor axis)} \\ &= \frac{2M_t}{\pi a b^2} \end{aligned}$$

### Bending of beams



$$\sigma_{zz} = E \epsilon_{zz} = \frac{E x}{R}$$

$$\epsilon_{zz} = \frac{x}{R}$$

all other  $\sigma = 0$

$$M = \int \sigma_{zz} x dA = \int \frac{E x^2}{R} dA = \frac{E I}{R}$$

$$\epsilon_{zz} = \frac{\sigma_{zz}}{E} = \frac{x}{R} = \frac{\partial u_z}{\partial z}$$

$$\epsilon_{xx} = \epsilon_{yy} = -\frac{v x}{R} = \frac{\partial u_x}{\partial x} = \frac{\partial u_y}{\partial y}$$

$$\frac{\partial u_z}{\partial x} + \frac{\partial u_x}{\partial z} = 0 = \gamma_{zx}$$

$$u_z = \frac{x z}{R} + w_0(x, y)$$

$$\gamma_{zx} = 0 = \frac{z}{R} + \frac{\partial w_0}{\partial x} + \frac{\partial u_x}{\partial z}$$

$$\frac{\partial u_x}{\partial z} = -\frac{z}{R} - \frac{\partial w_0}{\partial x}$$

$$-\frac{\partial u_z}{\partial y} = \frac{\partial u_y}{\partial z} = -\frac{\partial w_0}{\partial y}$$

(57)

$$u_x = -\frac{z^2}{2R} - z \frac{\partial w_0}{\partial x} + u_0(x, y)$$

$$u_y = -z \frac{\partial w_0}{\partial y} + v_0(x, y)$$

$$\left. \begin{aligned} \frac{\partial u_x}{\partial x} &= -z \frac{\partial^2 w_0}{\partial x^2} + \frac{\partial u_0}{\partial x} = -\frac{v_x}{R} \\ \frac{\partial u_y}{\partial y} &= -z \frac{\partial^2 w_0}{\partial y^2} + \frac{\partial v_0}{\partial y} = -\frac{v_x}{R} \end{aligned} \right\} \begin{array}{l} \text{must be valid for} \\ \text{any } z \end{array}$$

$$\frac{\partial^2 w_0}{\partial x^2} = \frac{\partial^2 w_0}{\partial y^2} = 0$$

$$\frac{\partial u_0}{\partial x} = -\frac{v_x}{R}$$

$$u_0 = -\frac{v_x x^2}{2R} + f_1(y)$$

$$v_0 = -\frac{v_x y}{R} + f_2(x)$$

$$u_x = -\frac{z^2}{2R} - z \frac{\partial w_0}{\partial x} - \frac{v_x x^2}{2R} + f_1(y)$$

$$u_y = -z \frac{\partial w_0}{\partial y} - \frac{v_x y}{R} + f_2(x)$$

recall  $\frac{\partial u_x}{\partial y} + \frac{\partial u_y}{\partial x} = 0$

$$-z \frac{\partial^2 w_0}{\partial x \partial y} + \frac{\partial f_1}{\partial y} - \frac{v_x y}{R} + \frac{\partial f_2}{\partial x} - z \frac{\partial^2 w_0}{\partial y \partial x} = 0$$

↑ ↑  
independent of z

$$\frac{\partial^2 w_0}{\partial x \partial y} = 0$$

collecting like terms  $\left\{ \begin{array}{l} \frac{\partial f_2}{\partial x} = -\alpha \\ \frac{\partial f_1}{\partial y} - \frac{v_x y}{R} = +\alpha \end{array} \right.$

$$f_2 = -\alpha x + \beta$$

$$f_1 = \alpha y + \frac{\nu y^2}{2R} + \gamma$$

$$\frac{\partial^2 w_0}{\partial x^2} = \frac{\partial^2 w_0}{\partial y^2} = 0 \quad (\text{from earlier})$$

$$w_0 = mx + ny + p$$

now,

$$u_x = -\frac{z^2}{2R} - zm - \frac{\nu x^2}{2R} + \alpha y + \frac{\nu y^2}{2R} + \gamma$$

$$u_y = -zn - \frac{\nu xy}{R} - \alpha x + \beta$$

find constants  $m, n, \alpha, \beta, \gamma$

$$\text{at } A \quad u_x = u_y = 0$$

$$\frac{\partial u_x}{\partial z} = \frac{\partial u_y}{\partial z} = \frac{\partial u_y}{\partial x} = 0$$

this sets all constants to zero

$$u_x = -\frac{z^2}{2R} - \frac{\nu x^2}{2R} + \frac{\nu y^2}{2R}$$

$$(u_x = -\frac{1}{2R}(z^2 + \nu(x^2 - y^2))) \quad \text{Timoshenko}$$

$$u_y = -\frac{\nu xy}{R}$$

deflection at  $y=x=0$

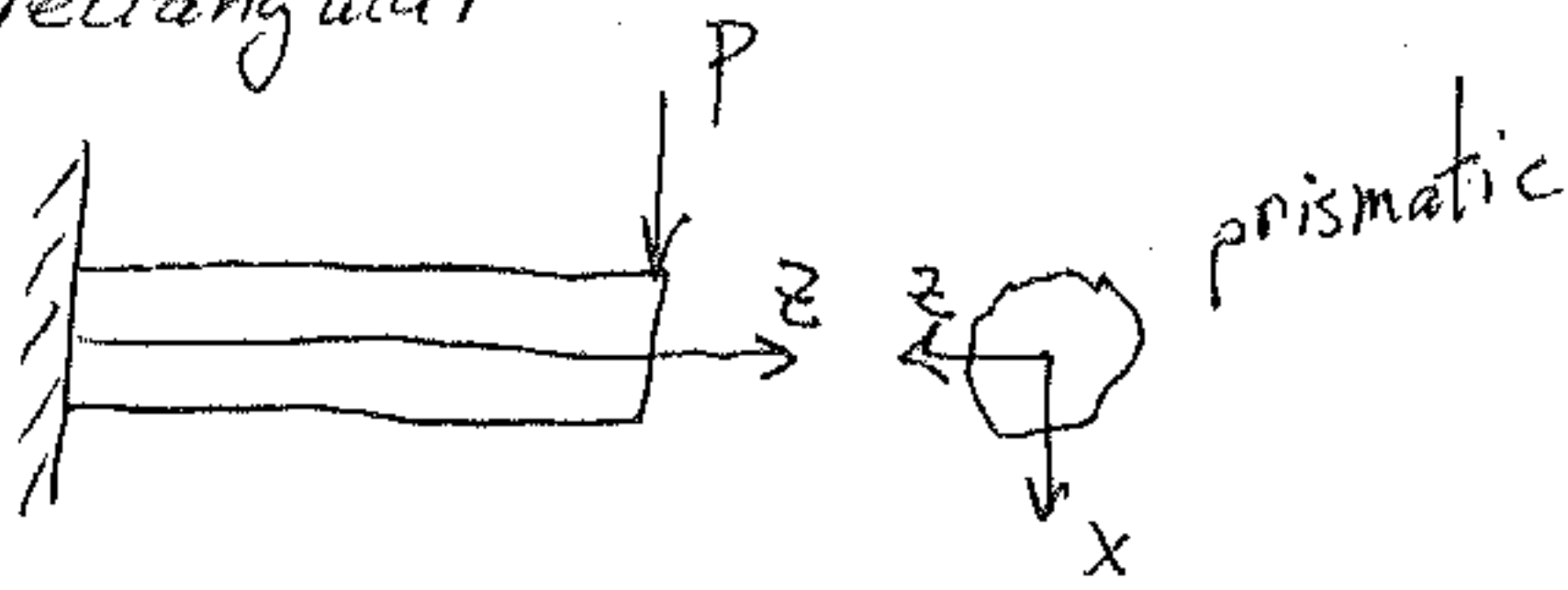
$$u_x|_{x=y=0} = -\frac{z^2}{2R} \quad \text{deflection curve}$$

Do plane sections remain plane?

$$u_z @ z=c \Rightarrow \frac{cx}{R} \quad \text{plane has rotated but consistent along } z$$

$$u_z = \frac{xz}{R} + mx + ny + p = \left(\frac{z}{R} + m\right)x + ny + p = \frac{zx}{R} \\ (m=n=p=0)$$

Bending of beams with other cross sections than rectangular



use St. Venant's semi-inverse method

$$\sigma_{zz} = -P(l-z)\frac{x}{I}$$

$$\tau_{xz} \neq 0$$

$$\tau_{yz} \neq 0$$

$$\sigma_{xx}, \sigma_{yy}, \tau_{xy} = 0$$

Equilibrium

$$\frac{\partial \tau_{xz}}{\partial z} = 0 \quad \frac{\partial \tau_{yz}}{\partial z} = 0$$

$$\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} = -\frac{Px}{I} = -\frac{\partial \sigma_{zz}}{\partial z}$$

B.C.'s

$$\tau_{xz} n_x + \tau_{yz} n_y = 0$$

$$\tau_{xz} \frac{dy}{ds} - \tau_{yz} \frac{dx}{ds} = 0$$

Compatibility

$$(1+\nu)\nabla^2 \sigma_{xx} + \sigma_{kk,xx} = 0$$

$$(1+\nu)\nabla^2 \sigma_{yy} + \sigma_{kk,yy} = 0$$

$$(1+\nu)\nabla^2 \sigma_{zz} + \sigma_{kk,zz} = 0$$

$$(1+\nu)\nabla^2 \sigma_{yz} + \sigma_{kk,yz} = 0$$

also for xz and xy



(60)

several components are zero

$$\left. \begin{aligned} \sigma_{zz,xx} &= 0 \\ \sigma_{zz,yy} &= 0 \\ (1+\nu)\nabla^2 \sigma_{zz} + \sigma_{zz,zz} &= 0 \end{aligned} \right\} \text{already satisfied by } \sigma_{zz} \text{ function}$$

$$\nabla^2 \sigma_{yz} = 0$$

$$\sigma_{zz,zx} + (1+\nu)\nabla^2 \sigma_{xz} = 0$$

$$\left. \begin{aligned} -\frac{P}{I} \frac{L}{(1+\nu)} &= \nabla^2 \sigma_{xz} \\ \nabla^2 \sigma_{yz} &= 0 \end{aligned} \right\} \text{must satisfy compatibility}$$

Introduce a stress function  $\phi(x,y)$ 

$$\tau_{xz} = \frac{\partial \phi}{\partial y} - \frac{Px^2}{2I} + f(y)$$

$$\tau_{yz} = -\frac{\partial \phi}{\partial x}$$

compatibility becomes

$$\frac{\partial}{\partial x} (\phi_{,xx} + \phi_{,yy}) = 0$$

$$\frac{\partial}{\partial y} (\phi_{,xx} + \phi_{,yy}) = \frac{\nu}{1+\nu} \frac{P}{I} - \frac{d^2 f}{dy^2}$$

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = \frac{\nu}{1+\nu} \frac{P}{I} - \frac{df}{dy} + C$$

(can be shown that  $C=0$ , assuming no twisting) - see TimoshenkoSubstitute  $\tau_{xz}, \tau_{yz}$  into B.C.'s

$$\frac{\partial \phi}{\partial y} \frac{dy}{ds} + \frac{\partial \phi}{\partial x} \frac{dx}{ds} = \frac{d\phi}{ds} = \left[ \frac{Px^2}{2I} - f(y) \right] \frac{dy}{ds}$$

if  $f(y) = \frac{Px^2}{2I}$  then  $\phi = \text{constant on boundary}$   
 $\phi = 0$  can be used